Cosmological constraints on the number of neutrino species

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Radiation content of the Universe

Microwave background

$$ho_{m{\gamma}} = rac{g_{m{\gamma}}}{(2\pi)^3} \, \int \mathrm{d}^3 q \; q \, f_{\mathrm{BE}}(q) = rac{\pi^2}{15} \, T_{m{\gamma}}^4$$

Neutrino background

$$ho_{m{
u}}^{
m act} = 3 \cdot rac{g_{m{
u}}}{(2\pi)^3} \, \int {
m d}^3 q \; q \, f_{m{
u}}(q) = N_{
m eff}^{
m act} \cdot rac{7\pi^2}{120} \, \left(rac{4}{11}
ight)^{4/3} T_{m{\gamma}}^4$$

standard model expectation:

$$N_{
m eff}^{
m act}=3.046$$
 [Mangano et al. (2008

Other light stuff?

$$ho_X = N_X \cdot rac{7\pi^2}{120} \, \left(rac{4}{11}
ight)^{4/3} T_{m \gamma}^4$$

Radiation content of the Universe

Putting it all together:



A few remarks on $N_{ m eff}$

is not a constant, in general

- increase through light decay products of massive particle
- decrease when particles go non-relativistic
- can be < 3.046, if neutrinos out of equilibrium; e.g., low reheating temperature:

[Ichikawa, Kawasaki, Takahashi (2005)]

- Cosmology can probe $N_{
 m eff}$
 - at BBN
 - at decoupling



$N_{ m eff}$ during BBN

- BBN element abundances depend on:
 - nuclear interaction rates
 - expansion rate



N_{eff} during BBN

- BBN element abundances depend on:
 - nuclear interaction rates
 - expansion rate



- Most sensitive probe:
 ⁴He abundance Y_p
- Y_p measured in metal-poor H-II regions
 - → subject to systematics



$N_{\rm eff}$ from measurements of Y **Recent measurements:** • $Y_{p} = 0.2561 \pm 0.0108$ [Aver, Olive, Skillman (2010)] Y_n = 0.2565 ± 0.001 (stat) ± 0.005 (syst) [Izotov, Thuan (2010)] $N_X + 3 = 3.68 + 0.80_{-0.70}$ $N_X + 3 = 3.65 + 1.97_{-1.57}$ (2σ) a) 4.5 $N_X + 3$ 3.5 =885.4±0.9s = 0.640524 $(H)_{=}(2.76\pm0.25)\times10^{-5}$ $(^{7}\text{Li}/\text{H})_{n} = (5.25 \pm 0.20) \times 10^{-10}$ 2.5 10 η_{10}

$N_{\rm eff}$ from CMB data

Matter-radiation equality

$$1 + z_{
m eq} = rac{\Omega_{
m m}}{\Omega_{
m r}} \simeq rac{\Omega_{
m m}h^2}{\Omega_{\gamma}h^2}rac{1}{1 + 0.2271N_{
m eff}}$$





$N_{\rm eff}$ from CMB data

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serious degeneracy with matter density!

- Anisotropic stress
 - present for neutrinos or other non-interacting particles, but not for photons
 - can help break degeneracy with matter density





- Iower limit from CMB alone (--> anisotropic stress)
- upper limit by combining with other data sets sensitive to matter density

[Komatsu et al. (2008)]

CMB+X bounds on $N_{ m eff}$

95%-credible intervals



- Precise numbers depend on cosmological model and data sets used
- ◆ Recent analysis: N_{eff} = 4.47 ^{+1.82}_{-1.74}
 CMB + SDSS-DR7-BAO + HST
 ∧CDM + neutrino mass + N_{eff}
 [JH, Hannestad, Lesgourgues, Rampf, Wong (2010)]

today: arXiv:1009.0866 including ACT small scale CMB data $N_{
m eff}$ = 4.56 ± 1.5 (95%)

CMB+X bounds on $N_{ m eff}$

95%-credible intervals



 Precise numbers depend on cosmological model and data sets used

Recent analysis: $N_{\rm eff}$ = 4.47 ^{+1.82} CMB + SDSS-DR7-BAO + HST ACDM + neutrino mass + $N_{\rm eff}$ [JH, Hannestad, Lesgourgues, Rampf, Wong (2010)]

PLANCK expected sensitivity $\sigma_{N_{\rm eff}} \approx 0.15$

[JH, Lesgourgues, Mangano (2007)]

Other hints: sterile neutrinos?



LSND & MiniBooNE anomalies may be resolved with CPviolating sterile neutrino oscillations in a "3+2" model (however: disfavoured by disappearance experiments)

[Karagiorgi, talk @ Neutrino 2010]

Sterile neutrino scenario



Sterile neutrino scenario

♦ 3+2: m_s < 0.45 eV (@ 95% c.l.)</p>





[JH, Hannestad, Raffelt, Tamborra, Wong (2010)]

Conclusions

- Cosmological observations are a powerful probe of the Universe's radiation content
- Compelling indirect evidence for the existence of a cosmic neutrino background, from both BBN and CMB+LSS data
- Two sterile neutrino species still compatible with cosmological data
- Planck will settle the issue!

