

# Short-BaseLine Electron Neutrino Disappearance

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Statistical Significance of the Gallium Anomaly  
[arXiv:1006.3244](https://arxiv.org/abs/1006.3244)

Hint of CPT Violation in Short-Baseline Electron Neutrino Disappearance  
[arXiv:1008.4750](https://arxiv.org/abs/1008.4750)

## Standard Model: Massless Neutrinos

- Standard Model:  $\nu_L, \nu_R^c \quad \Rightarrow \quad$  no Dirac mass term

$$\mathcal{L}^D = m^D (\overline{\nu}_L \nu_R + \overline{\nu}_R \nu_L)$$

- Majorana Neutrino:  $\nu^c = \nu$

- $\nu_R^c = \nu_R \quad \Rightarrow \quad$  Majorana mass term

$$\mathcal{L}^M = \frac{1}{2} m^M (\overline{\nu}_L \nu_R^c + \overline{\nu}_R^c \nu_L)$$

- Standard Model: Majorana mass term **not** allowed by  $SU(2)_L \times U(1)_Y$

(no Higgs triplet)

- Neutrinos are special in the Standard Model: the only **neutral fermions**
- In extensions of SM neutrinos can mix with non-SM fermions

$$L_{\alpha L} = \begin{pmatrix} \nu_{\alpha L} \\ \alpha_L \end{pmatrix} \quad \tilde{\Phi} = i\sigma_2 \Phi^* = \begin{pmatrix} \phi^0 \\ \phi^- \end{pmatrix} \xrightarrow[\text{Breaking}] {\text{Symmetry}} \begin{pmatrix} v/\sqrt{2} \\ 0 \end{pmatrix}$$

$(\alpha = e, \mu, \tau)$

$\overline{L_{\alpha L}} \tilde{\Phi}$  can be coupled to new non-SM chiral fermion fields  $f_{\beta R}$

Dirac mass terms  $\sim \overline{L_{\alpha L}} \tilde{\Phi} f_{\beta R}$  + Majorana mass terms  $\sim \overline{f_{\beta R}^C} f_{\beta R}$

$f_{\beta R}$  are often called **Right-Handed Neutrinos**:  $f_{\beta R} \rightarrow \nu_{\beta R}$

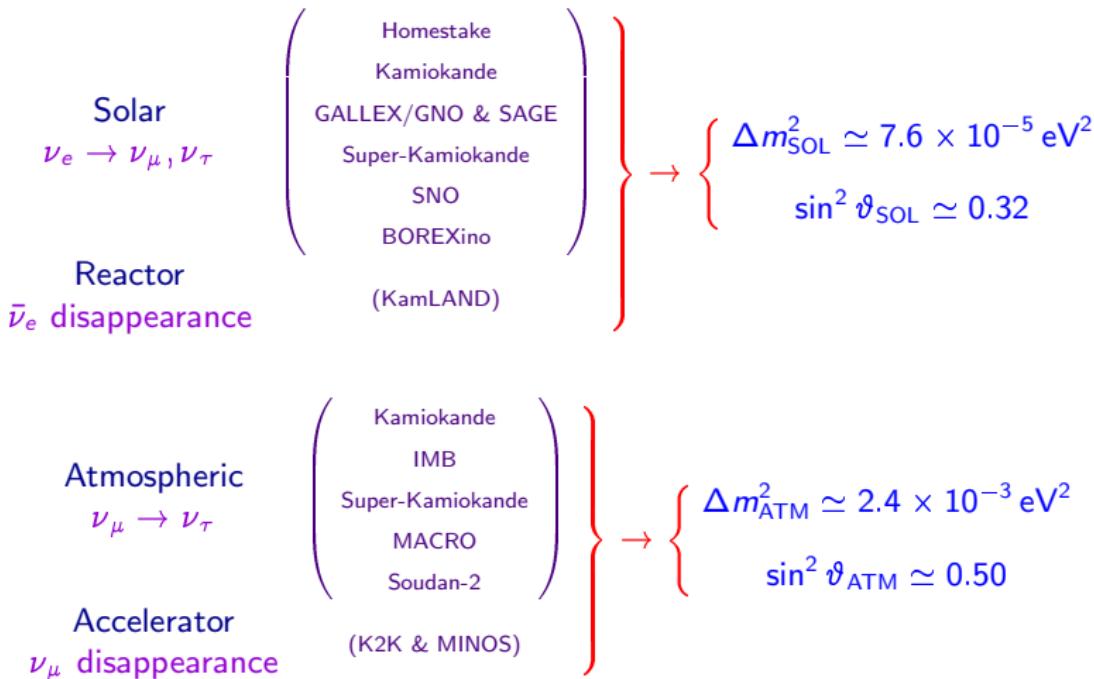
- If  $f_{\beta R}$  are light, they are called **Sterile Neutrinos**:

$$\nu_{s_{\beta} L} = f_{\beta R}^C$$

# Sterile Neutrinos

- ▶ Sterile means No Standard Model Interactions
- ▶ Obviously no electromagnetic interactions as normal active neutrinos
- ▶ Thus Sterile means No Standard Weak Interactions
- ▶ But Sterile Neutrinos are not absolutely sterile:
  - ▶ Gravitational Interactions
  - ▶ New Non-Standard Interactions of the Physics Beyond the Standard Model which generates the masses of sterile neutrinos
- ▶ Extremely interesting and powerful window on Physics Beyond the SM
- ▶ Active Neutrinos ( $\nu_e, \nu_\mu, \nu_\tau$ ) can oscillate into Sterile Neutrinos ( $\nu_s$ )
- ▶ Observables:
  - ▶ disappearance of Active Neutrinos
  - ▶ indirect evidence through combined fit of data

# Solar and Atmospheric Neutrino Oscillations



Two scales of  $\Delta m^2 \iff$  Three-Neutrino Mixing

$$\Delta m_{\text{SOL}}^2 = \Delta m_{21}^2 \simeq 7.6 \times 10^{-5} \text{ eV}^2$$

$$\Delta m_{\text{ATM}}^2 \simeq |\Delta m_{31}^2| \simeq |\Delta m_{32}^2| \simeq 2.4 \times 10^{-3} \text{ eV}^2$$

- ▶ New Short-BaseLine Oscillations:  $\frac{L}{E} \lesssim 1 \frac{\text{m}}{\text{MeV}} \implies \Delta m_{\text{SBL}}^2 \gtrsim 1 \text{ eV}^2$
- ▶ Necessary introduction of at least one new massive neutrino: 4- $\nu$  Mixing

$$\Delta m_{\text{SBL}}^2 = \Delta m_{41}^2$$

Mass Basis:  $\nu_1 \quad \nu_2 \quad \nu_3 \quad \nu_4$

Flavor Basis:  $\nu_e \quad \nu_\mu \quad \nu_\tau \quad \nu_s$

- ▶ Effective SBL Oscillation Probabilities:

$$▶ P_{\nu_\alpha \rightarrow \nu_\beta}^{\text{SBL}} = \sin^2 2\vartheta_{\alpha\beta} \sin^2 \left( \frac{\Delta m_{\text{SBL}}^2 L}{4E} \right) \quad (\alpha \neq \beta)$$

$$▶ P_{\nu_\alpha \rightarrow \nu_\alpha}^{\text{SBL}} = 1 - \sin^2 2\vartheta_{\alpha\alpha} \sin^2 \left( \frac{\Delta m_{\text{SBL}}^2 L}{4E} \right)$$

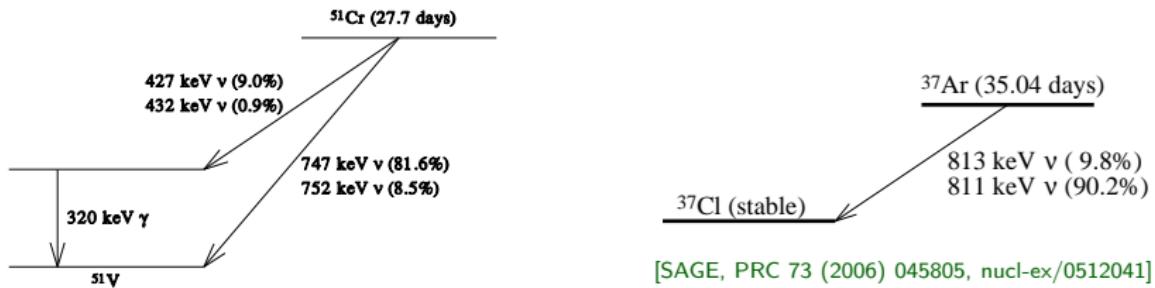
# Gallium Anomaly

## Gallium Radioactive Source Experiments

Tests of the solar neutrino detectors **GALLEX** (Cr1, Cr2) and **SAGE** (Cr, Ar)

Detection Process:  $\nu_e + {}^{71}\text{Ga} \rightarrow {}^{71}\text{Ge} + e^-$

$\nu_e$  Sources:  $e^- + {}^{51}\text{Cr} \rightarrow {}^{51}\text{V} + \nu_e$        $e^- + {}^{37}\text{Ar} \rightarrow {}^{37}\text{Cl} + \nu_e$



[SAGE, PRC 73 (2006) 045805, nucl-ex/0512041]

[SAGE, PRC 59 (1999) 2246, hep-ph/9803418]

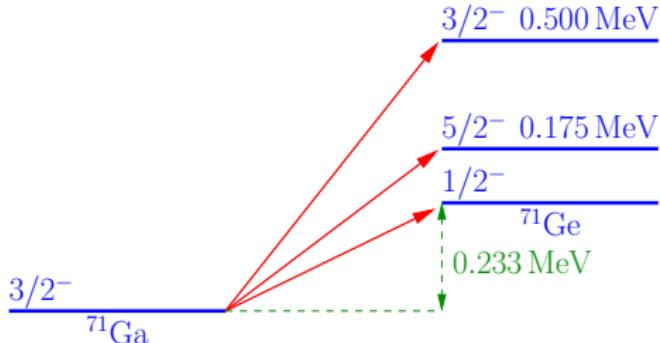
$$R_{\text{Ga}} = 0.76^{+0.09}_{-0.08}$$

[Giunti, Laveder, arXiv:1006.3244]

Haxton cross section and uncertainty

[Haxton, PLB 431 (1998) 110, nucl-th/9804011]

- Deficit could be due to overestimate of  
 $\sigma(\nu_e + {}^{71}\text{Ga} \rightarrow {}^{71}\text{Ge} + e^-)$
- Calculation: Bahcall, PRC 56 (1997) 3391, hep-ph/9710491



- $\sigma_{\text{G.S.}}$  related to measured  $\sigma(e^- + {}^{71}\text{Ge} \rightarrow {}^{71}\text{Ga} + \nu_e)$ :

$$\sigma_{\text{G.S.}}({}^{51}\text{Cr}) = 55.3 \times 10^{-46} \text{ cm}^2 (1 \pm 0.004)_{3\sigma}$$

$$\sigma({}^{51}\text{Cr}) = \sigma_{\text{G.S.}}({}^{51}\text{Cr}) \left( 1 + 0.669 \frac{\text{BGT}_{175 \text{ keV}}}{\text{BGT}_{\text{G.S.}}} + 0.220 \frac{\text{BGT}_{500 \text{ keV}}}{\text{BGT}_{\text{G.S.}}} \right)$$

- Contribution of Excited States only 5%!

► Bahcall:

[Bahcall, PRC 56 (1997) 3391, hep-ph/9710491]

from  $p + {}^{71}\text{Ga} \rightarrow {}^{71}\text{Ge} + n$  measurements [Krofcheck et al., PRL 55 (1985) 1051]

$$\frac{\text{BGT}_{175 \text{ keV}}}{\text{BGT}_{\text{G.S.}}} < 0.056 \Rightarrow \frac{\text{BGT}_{175 \text{ keV}}}{\text{BGT}_{\text{G.S.}}} = \frac{0.056}{2} \quad \frac{\text{BGT}_{500 \text{ keV}}}{\text{BGT}_{\text{G.S.}}} = 0.146$$

$3\sigma$  lower limit:  $\frac{\text{BGT}_{175 \text{ keV}}}{\text{BGT}_{\text{G.S.}}} = \frac{\text{BGT}_{500 \text{ keV}}}{\text{BGT}_{\text{G.S.}}} = 0$

$3\sigma$  upper limit:  $\frac{\text{BGT}_{175 \text{ keV}}}{\text{BGT}_{\text{G.S.}}} < 0.056 \times 2 \quad \frac{\text{BGT}_{500 \text{ keV}}}{\text{BGT}_{\text{G.S.}}} = 0.146 \times 2$

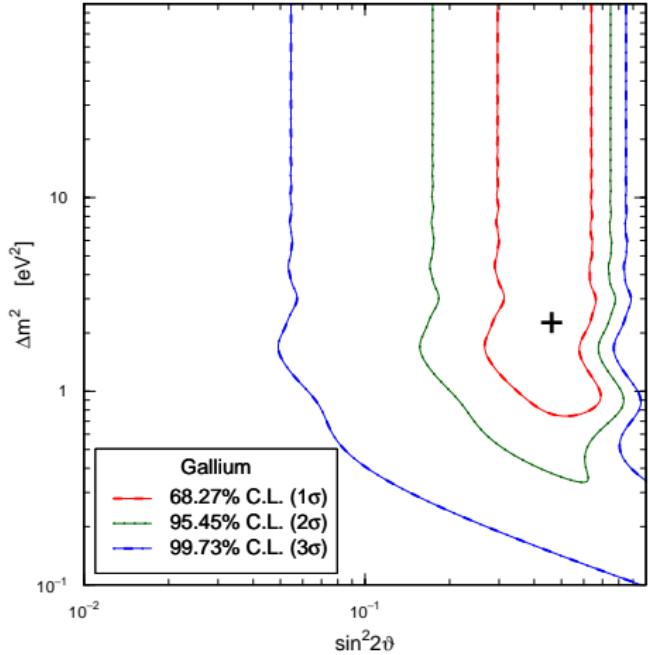
$$\sigma({}^{51}\text{Cr}) = 58.1 \times 10^{-46} \text{ cm}^2 \left(1^{+0.036}_{-0.028}\right)_{1\sigma}$$

► Haxton:

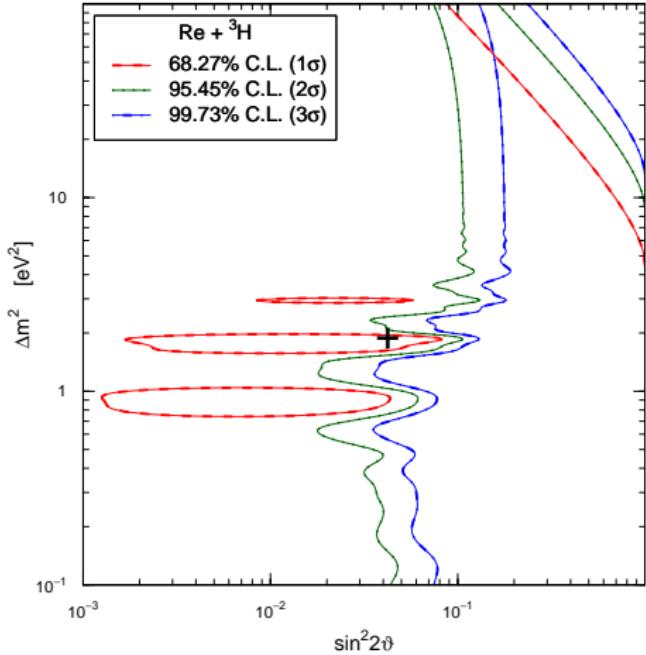
[Hata, Haxton, PLB 353 (1995) 422, nucl-th/9503017; Haxton, PLB 431 (1998) 110, nucl-th/9804011]

"a sophisticated shell model calculation is performed ... for the transition to the first excited state in  ${}^{71}\text{Ge}$ . The calculation predicts destructive interference between the  $(p, n)$  spin and spin-tensor matrix elements."

$$\sigma({}^{51}\text{Cr}) = 63.9 \times 10^{-46} \text{ cm}^2 (1 \pm 0.106)_{1\sigma}$$



[Giunti, Laveder, arXiv:1006.3244]

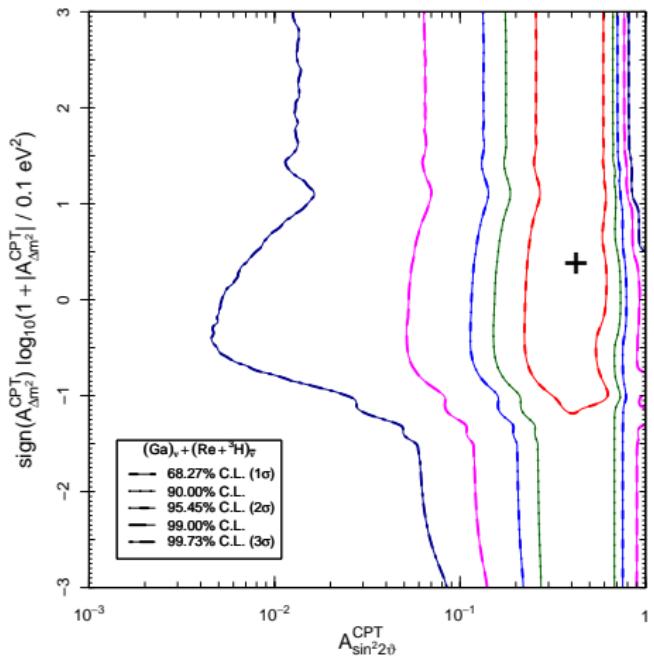


[Giunti, Laveder, arXiv:1005.4599]

$\Delta m_{\text{SBL}}^2 \gtrsim 1 \text{ eV}^2$  is OK

$\sin^2 2\theta_\nu > \sin^2 2\theta_{\bar{\nu}}$  CPT violation?

Parameter Goodness-Of-Fit:  $\Delta\chi^2_{\min} = 12.1$ , NDF = 2, GoF = 0.2%

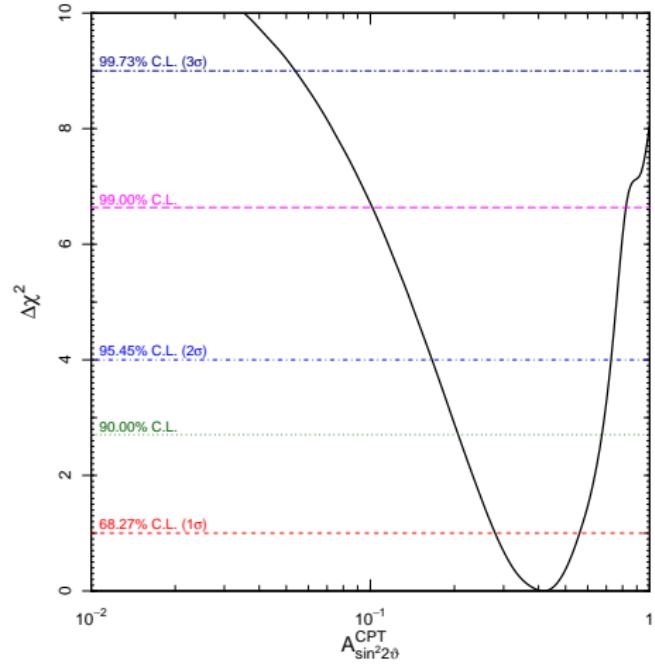


[Giunti, Laveder, arXiv:1008.4750]

$$A_{\sin^2 2\vartheta}^{\text{CPT}} = \sin^2 2\vartheta_\nu - \sin^2 2\vartheta_{\bar{\nu}}$$

$$A_{\Delta m^2}^{\text{CPT}} = \Delta m_\nu^2 - \Delta m_{\bar{\nu}}^2$$

$$A_{\sin^2 2\vartheta}^{\text{CPT}} > 0.055 \text{ at } 3\sigma$$



$$(A_{\sin^2 2\vartheta}^{\text{CPT}})_{\text{bf}} = 0.42$$

$$(A_{\Delta m^2}^{\text{CPT}})_{\text{bf}} = 0.37 \text{ eV}^2$$

$$A_{\sin^2 2\vartheta}^{\text{CPT}} > 0 \text{ at } 3.5\sigma.$$

## Future

- ▶ New Gallium source experiments:  $\nu_e$  disappearance [Gavrin et al, arXiv:1006.2103]
- ▶ CPT test:  $\nu_e$  and  $\bar{\nu}_e$  disappearance
- ▶ Beta-Beam experiments:

$$N(A, Z) \rightarrow N(A, Z + 1) + e^- + \bar{\nu}_e \quad (\beta^-)$$

$$N(A, Z) \rightarrow N(A, Z - 1) + e^+ + \nu_e \quad (\beta^+)$$

- ▶ Neutrino Factory experiments:

$$\mu^+ \rightarrow \bar{\nu}_\mu + e^+ + \nu_e$$

$$\mu^- \rightarrow \nu_\mu + e^- + \bar{\nu}_e$$

- ▶ New  $\nu_e$  and  $\bar{\nu}_e$  radioactive source experiments with low-threshold neutrino elastic scattering detectors? (as Borexino or liquid Argon TPC)

## Conclusions

- ▶ Gallium Anomaly may be a signal of Short-Baseline  $\nu_e$  disappearance with  $\Delta m^2 \gtrsim 1 \text{ eV}^2$  and  $\sin^2 2\vartheta \gtrsim 0.1$
- ▶ Tension with reactor  $\bar{\nu}_e$  disappearance limit  $\sin^2 2\vartheta \lesssim 0.1$
- ▶ Hint of CPT violation:  $A_{\sin^2 2\vartheta}^{\text{CPT}} > 0$  at  $3.5\sigma$ .
- ▶ Needed high-precision  $\nu_e$  and  $\bar{\nu}_e$  disappearance experiments
- ▶ Short-Baseline  $\nu_e$  disappearance maybe connected with LSND and MiniBooNE  $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$  signal (work in progress)