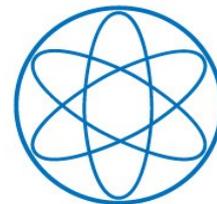
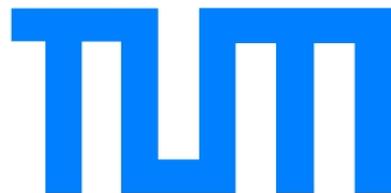


# Neutrino bounds on dark matter

Alejandro Ibarra

Technische Universität München



NOW 2012  
10 September 2012

# Introduction

Many pieces of evidence for particle dark matter. However, very little is known about the properties of the dark matter particle:

**Spin:** 0 or 1/2 or 1 or 3/2 (possibly higher if composite)

**Parity:** + or -

**Mass:**  $10^{-15}$  GeV  $\longrightarrow$   $10^{15}$  GeV  
(axions) (WIMPzillas)

**Self-annihilation cross section:**  $10^{-40}$  pb  $\longrightarrow$   $10^{-5}$  pb  
(gravitinos) (neutralinos)

**Interaction cross section with nucleons:**  $10^{-40}$  pb  $\longrightarrow$   $10^{-5}$  pb  
(gravitinos) (neutralinos)

**Lifetime:**  $10^9$  years  $\longrightarrow$  infinity

# Direct detection

DM nucleus  $\rightarrow$  DM nucleus



# Indirect detection

DM DM  $\rightarrow \gamma X, e^+ X, \nu X \dots$  (annihilation)

DM  $\rightarrow \gamma X, e^+ X, \nu X \dots$  (decay)

# Collider searches

pp  $\rightarrow$  DM X

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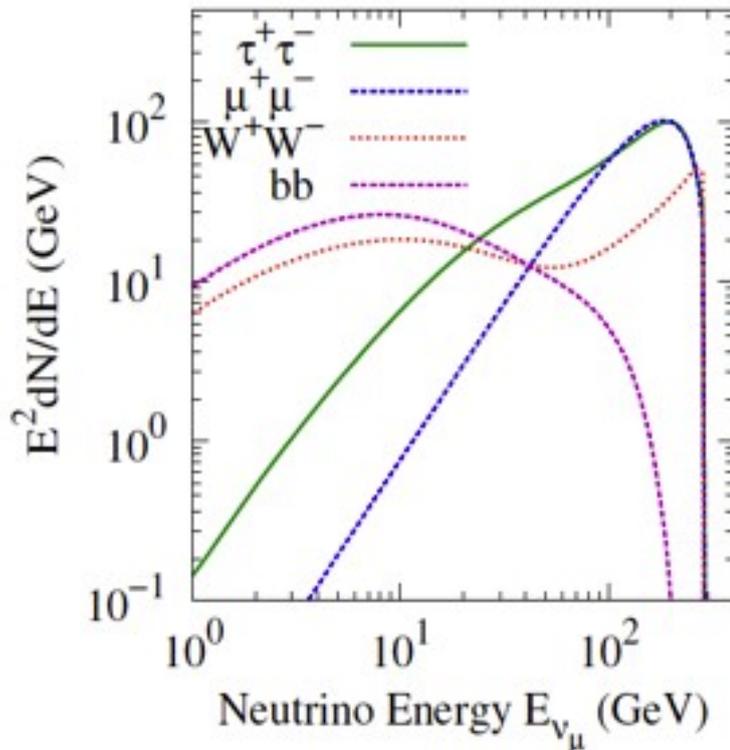
**2<sup>nd</sup> part**  $\longrightarrow$  **Interaction cross section with nucleons:**  $10^{-40}$  pb  $\longrightarrow$   $10^{-5}$  pb  
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# Limits on the annihilation cross-section

Neutrinos from dark matter annihilations in the Milky Way halo

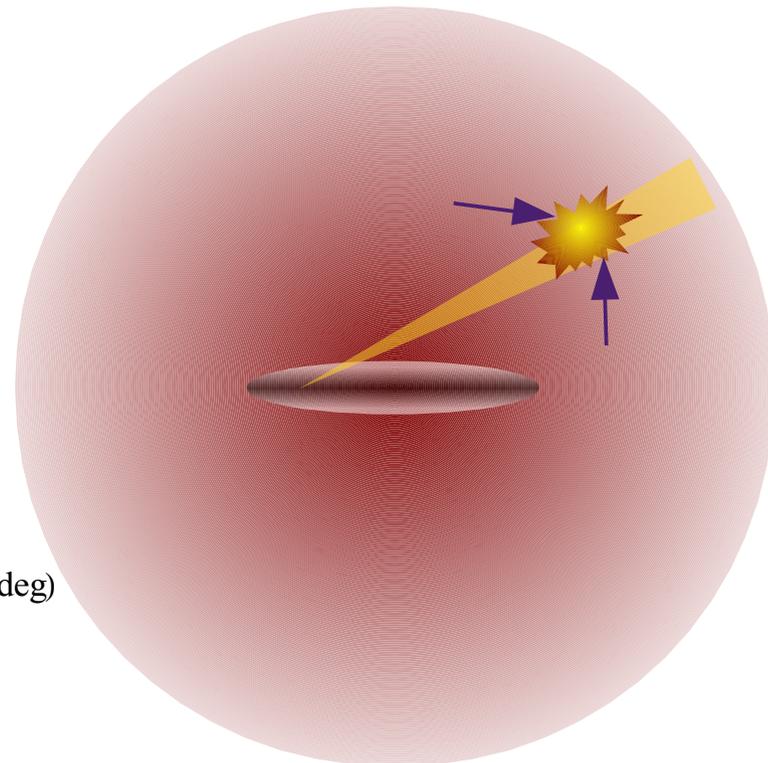
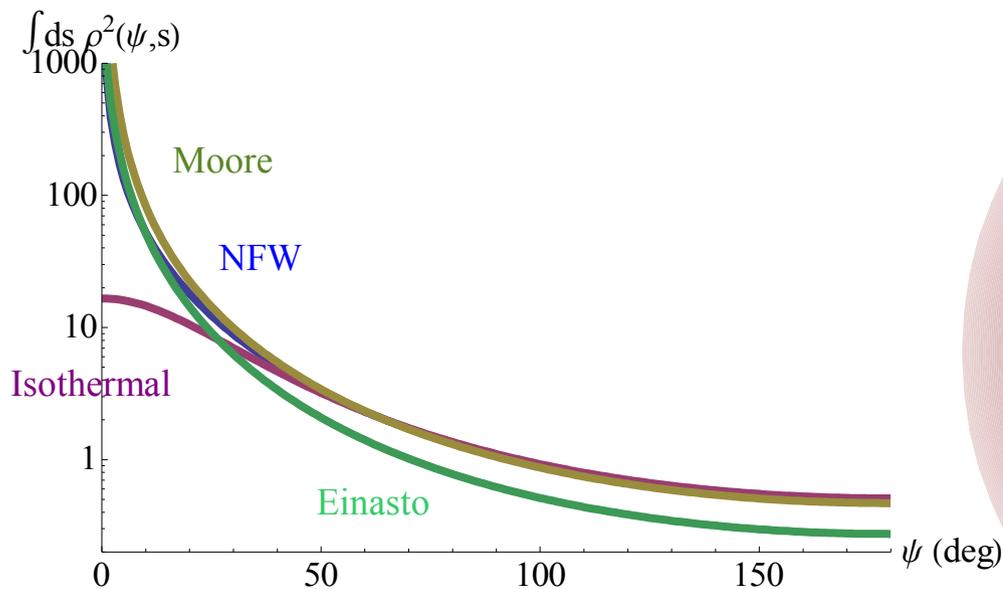
$$\frac{dJ_{\text{halo}}}{dE_\nu} = \frac{1}{4\pi} \underbrace{\left[ \frac{\langle \sigma_{\text{ann}} v \rangle}{2m_{\text{DM}}^2} \sum_f \frac{dN_\nu^f}{dE_\nu} B_f \right]}_{\text{Source term (particle physics)}} \times \underbrace{\int_{\text{l.o.s.}} \rho^2(\vec{l}) d\vec{l}}_{\text{Line-of-sight integral (astrophysics)}}$$



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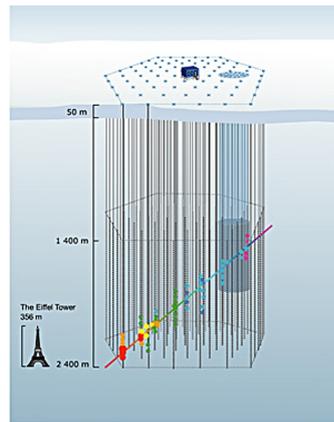


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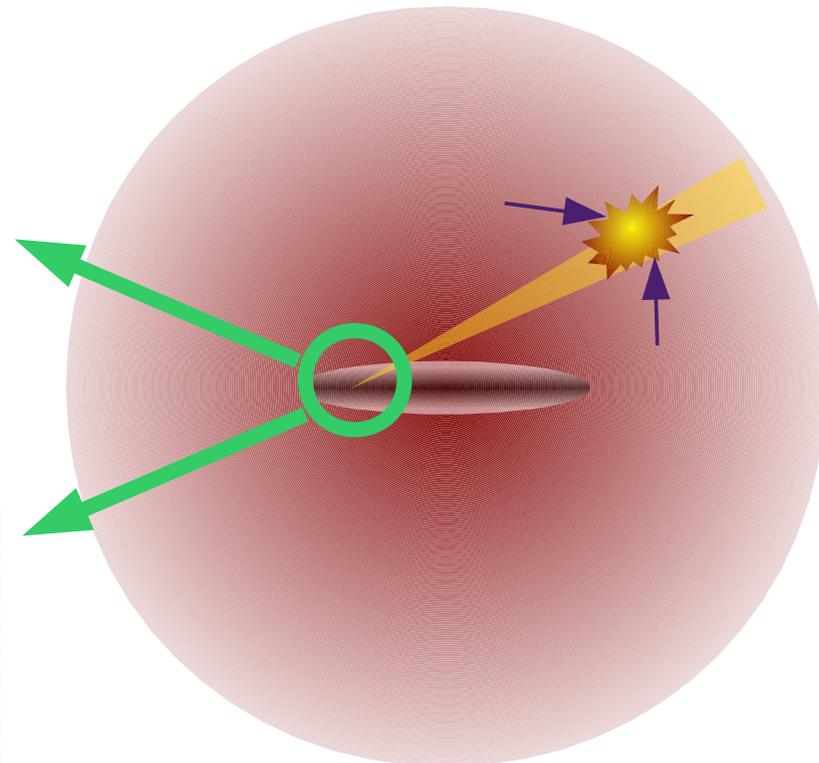
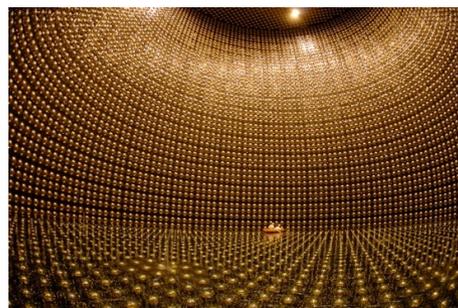
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IceCube

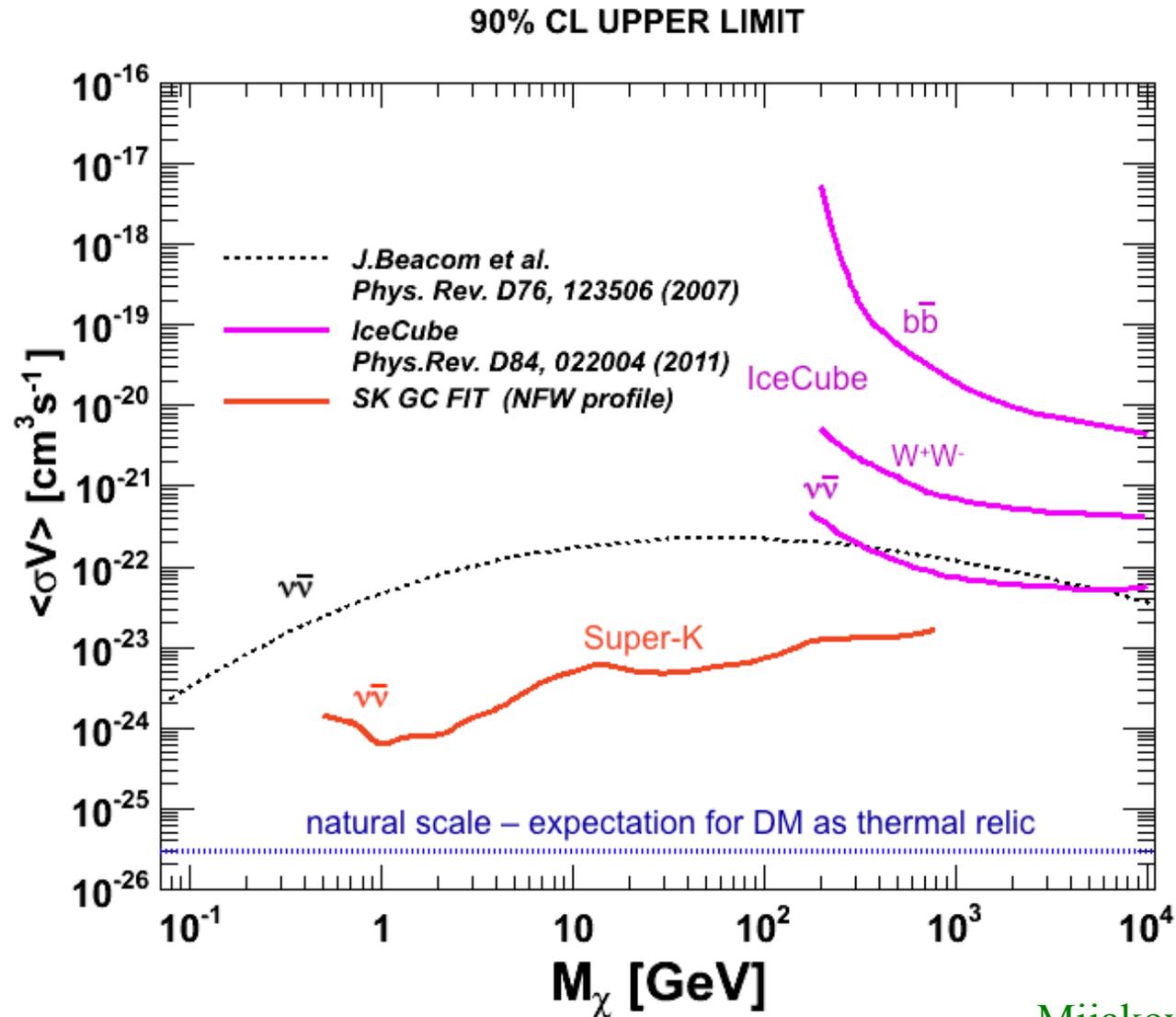


SuperK



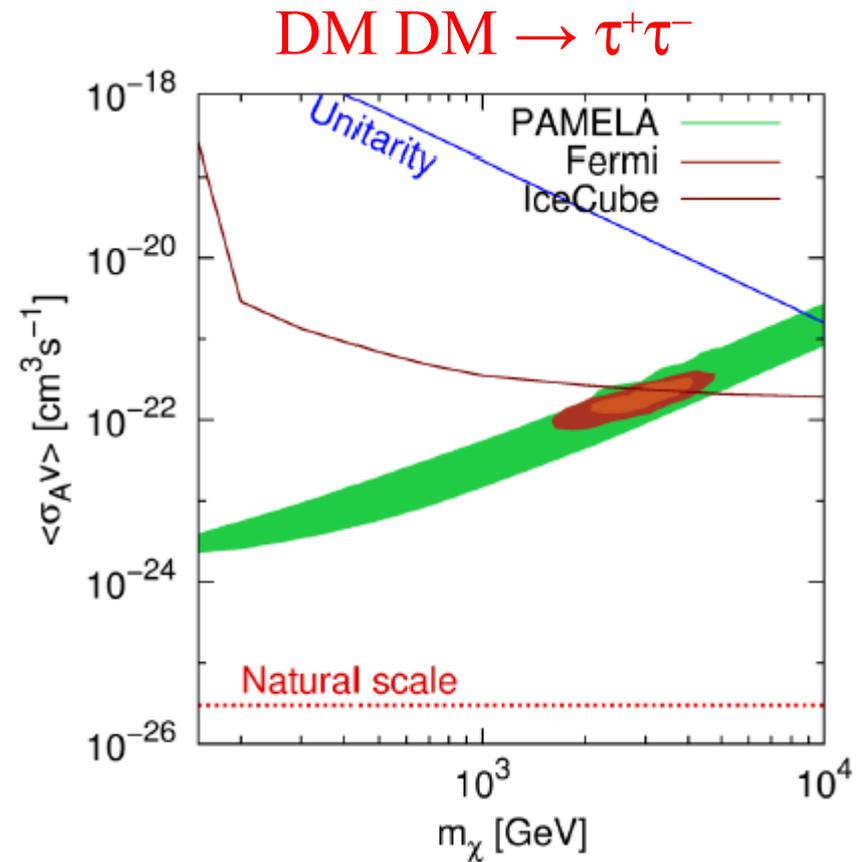
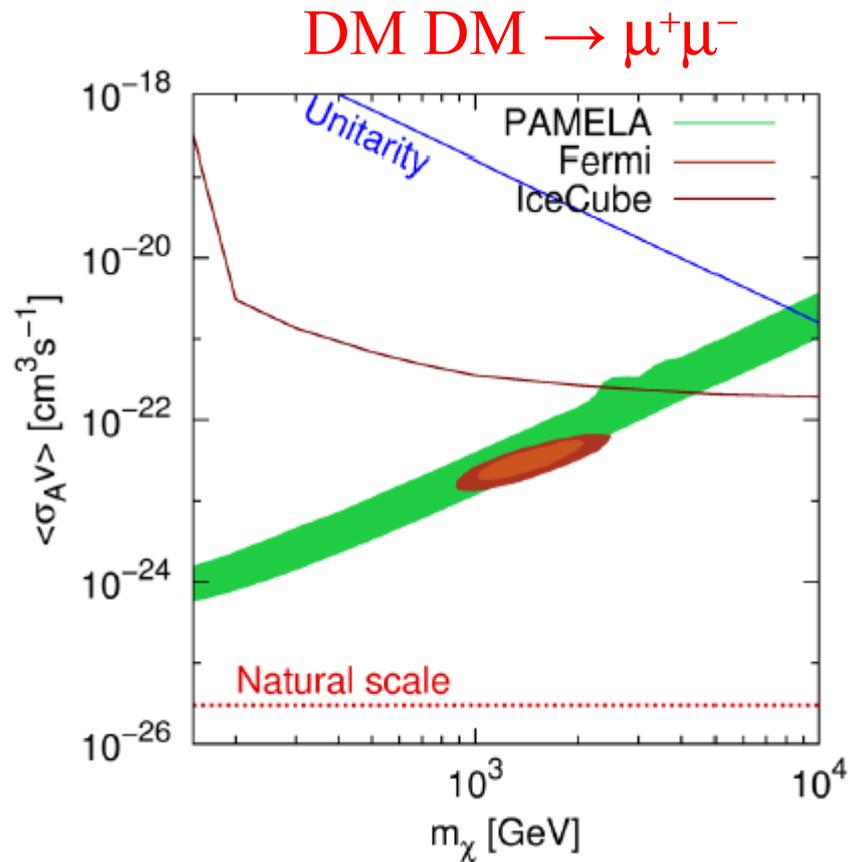
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# Limits on the annihilation cross-section

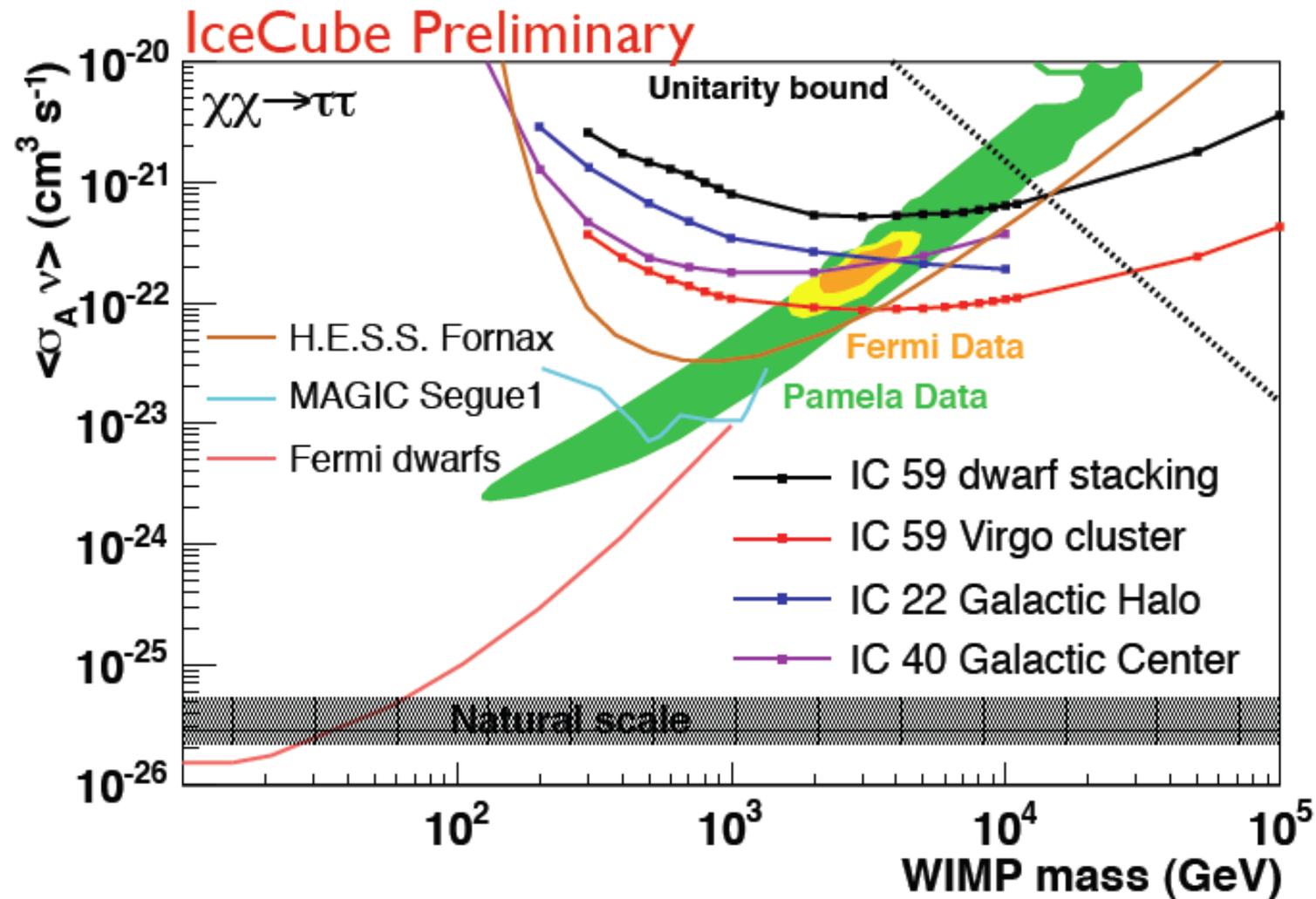
Neutrinos from dark matter annihilations in the Milky Way halo



Abbasi et al.  
arXiv:1101.3349

# Limits on the annihilation cross-section

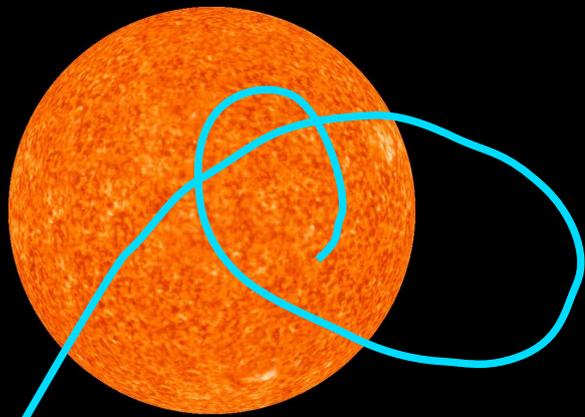
Neutrinos from dark matter annihilations in dwarf galaxies & galaxy clusters



Courtesy of Carsten Rott

# Limits on the scattering cross-section

- If the dark matter particles have a “sizable” interaction cross section with ordinary matter, they can be captured inside the Sun and inside the Earth.

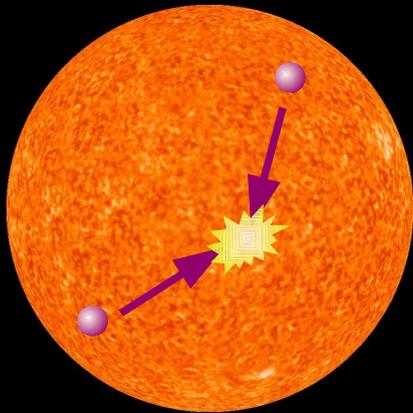


capture rate  $\propto \sigma_{\text{WIMP,p}}$



# Limits on the scattering cross-section

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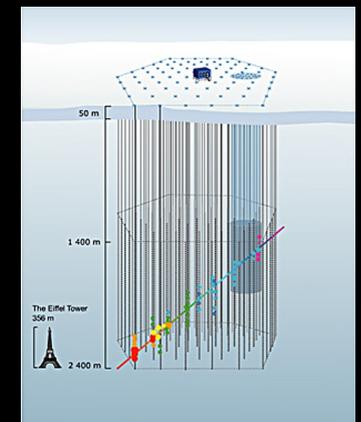
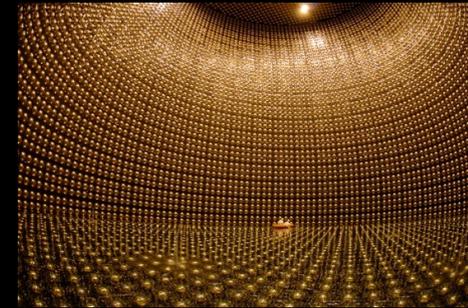
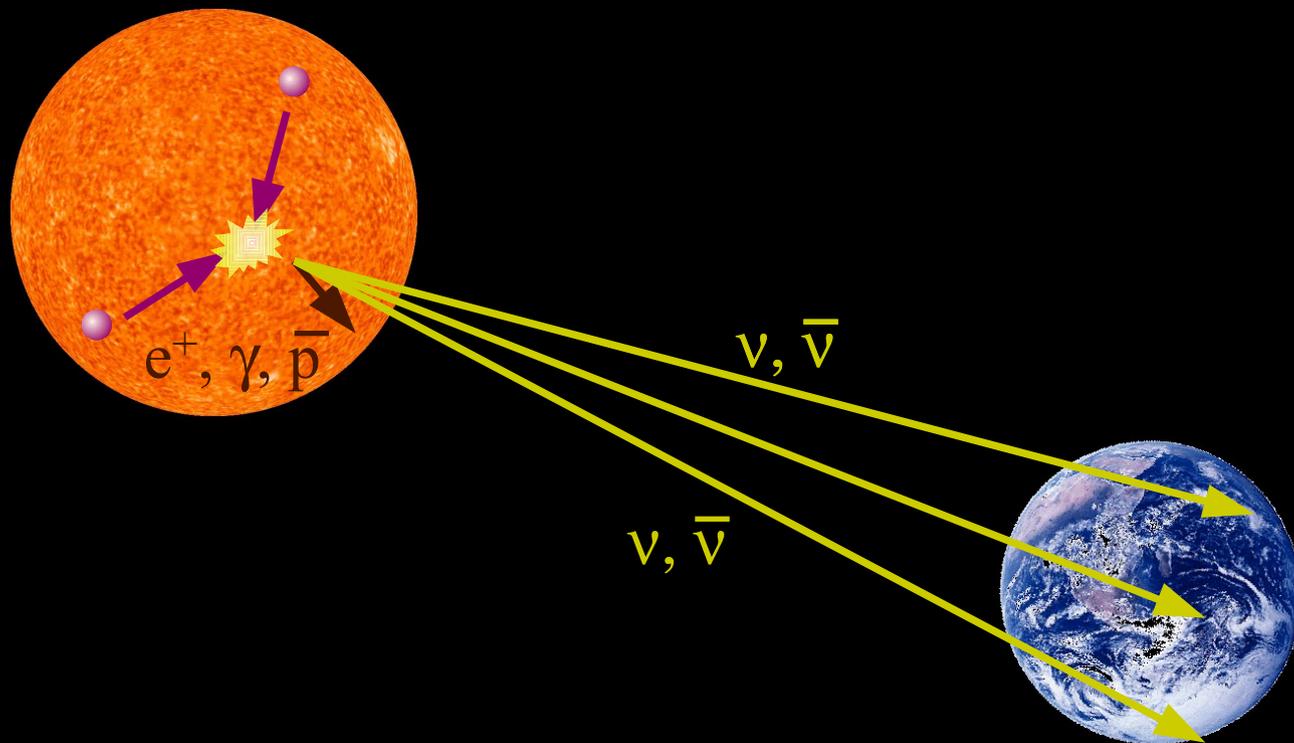
annihilation rate  $\propto \langle \sigma v \rangle$

$$\Gamma_A \simeq \frac{C_c}{2}$$



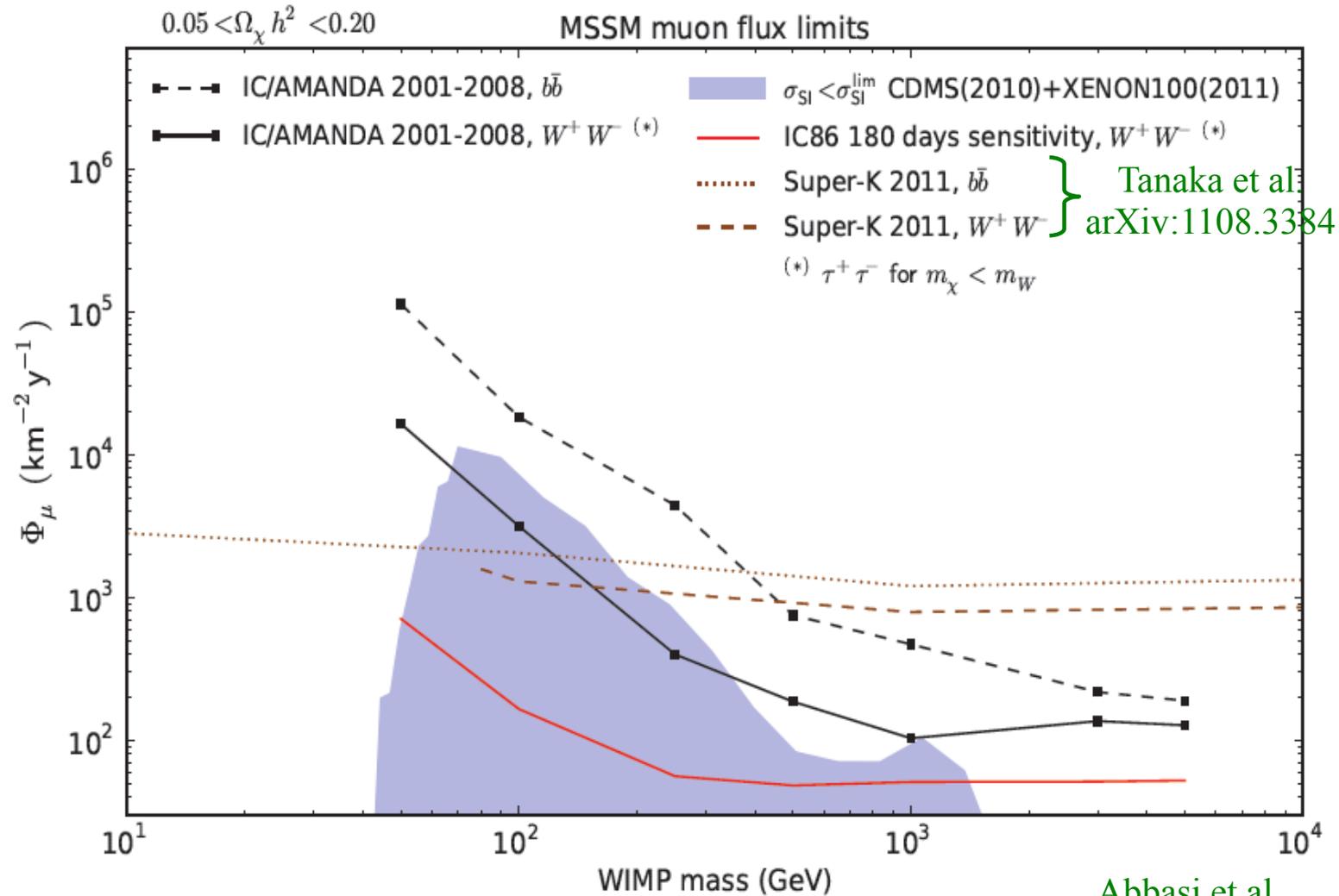
# Limits on the scattering cross-section

- If the dark matter particles have a “sizable” interaction cross section with ordinary matter, they can be captured inside the Sun and inside the Earth.
- DM particles captured inside the Sun/Earth can annihilate.
- The annihilation produces a neutrino flux which might be detected in neutrino observatories. All other annihilation products (gammas, positrons, antiprotons...) are absorbed before escaping the Sun/Earth.



# Limits on the scattering cross-section

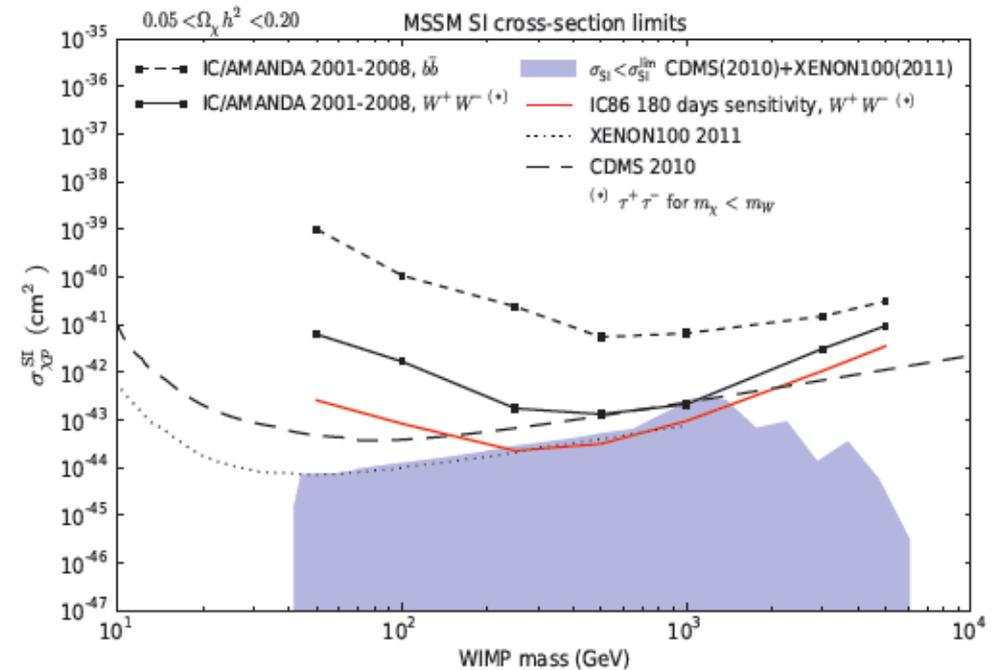
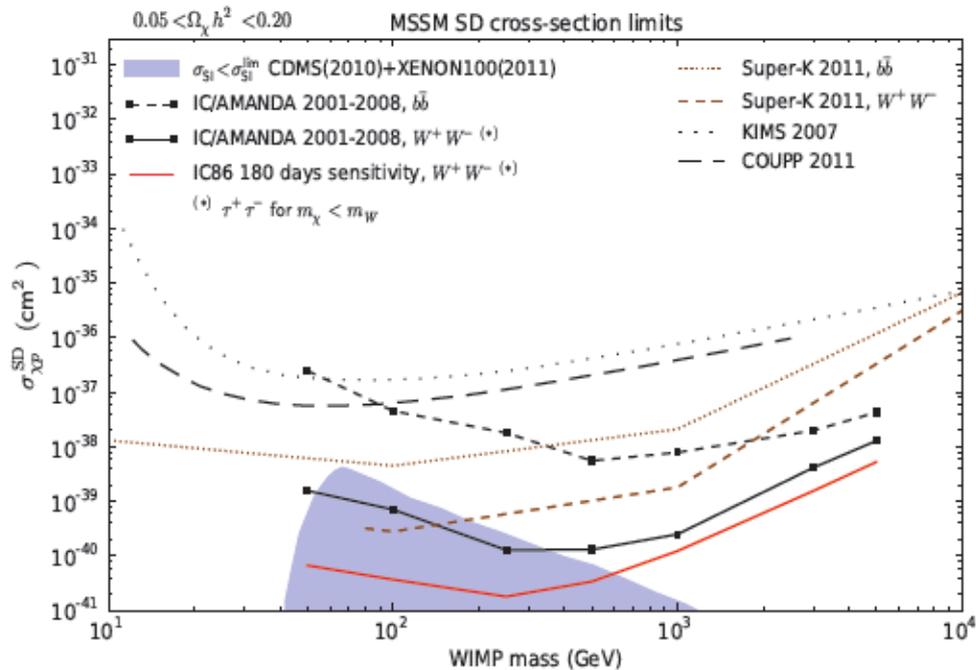
90% C.L. upper limits on the muon flux from neutralino annihilations in the Sun:



Abbasi et al.  
arXiv:1112.1840

# Limits on the scattering cross-section

## Limits on the SD and SI scattering cross sections

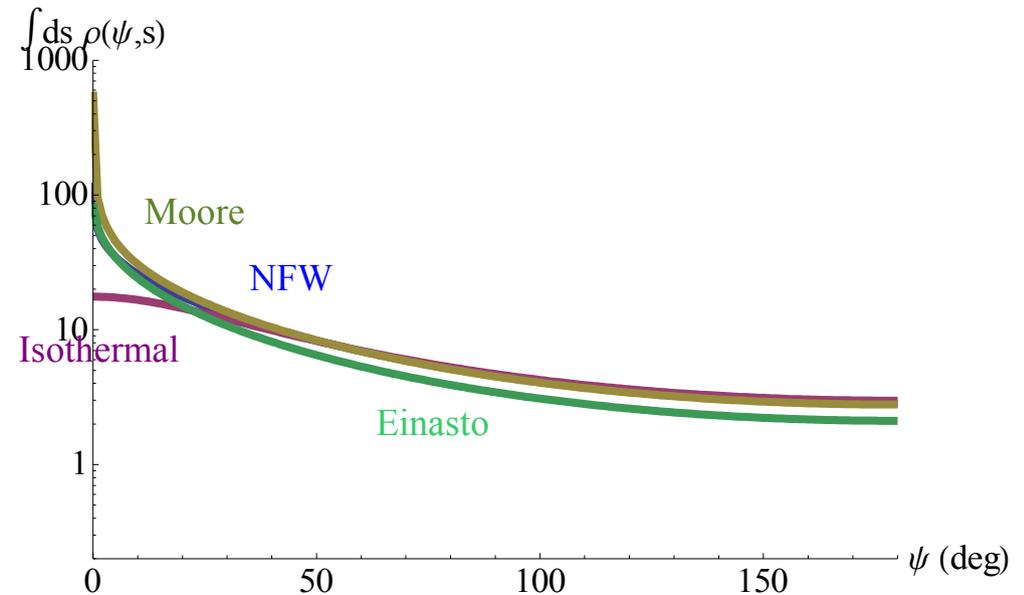
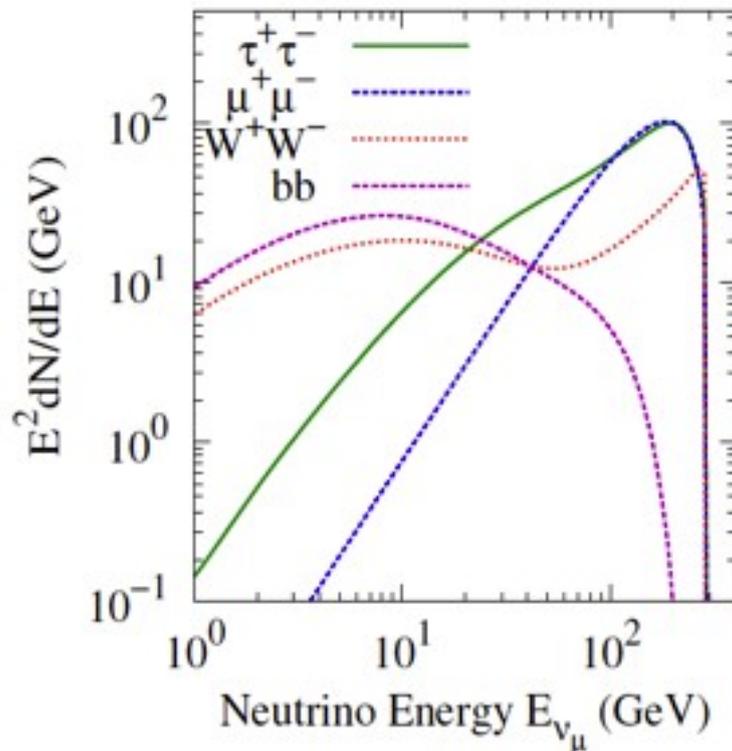


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# Limits on the dark matter lifetime

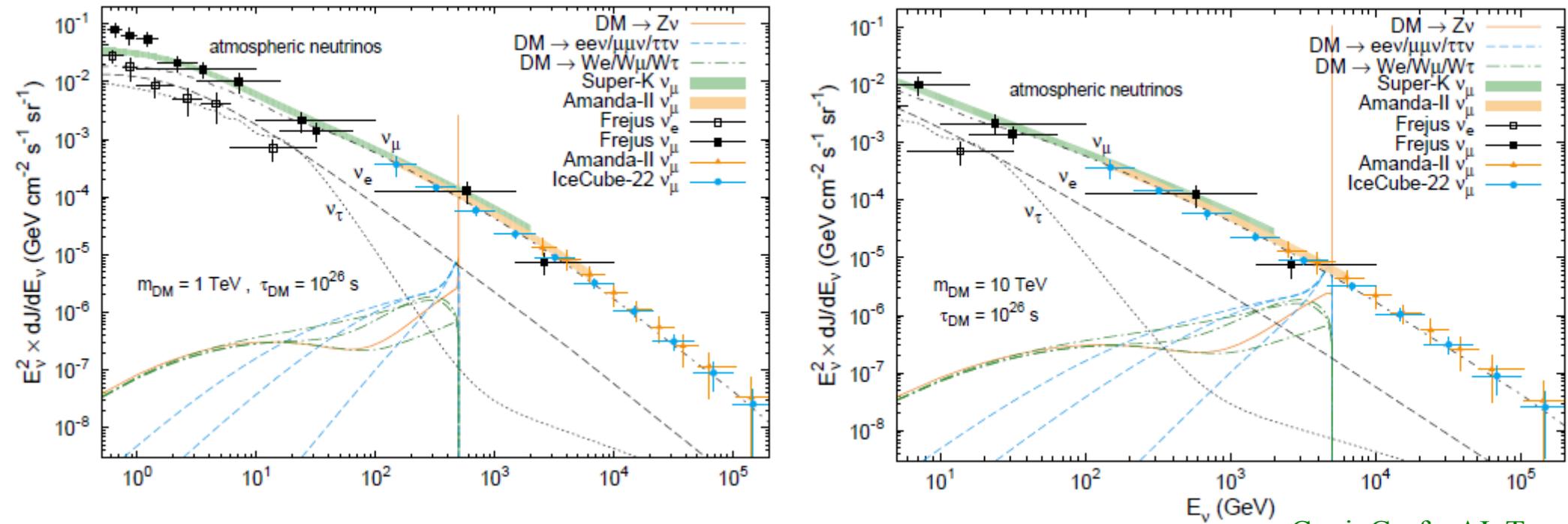
## Neutrinos from dark matter decay in the galactic halo

$$\frac{dJ_{\text{halo}}}{dE_\nu} = \frac{1}{4\pi} \underbrace{\left[ \frac{1}{\tau_{\text{DM}} m_{\text{DM}}} \sum_f \frac{dN_\nu^f}{dE_\nu} B_f \right]}_{\text{Source term (particle physics)}} \times \underbrace{\int_{\text{l.o.s.}} \rho(\vec{l}) d\vec{l}}_{\text{Line-of-sight integral (astrophysics)}}$$



# Limits on the dark matter lifetime

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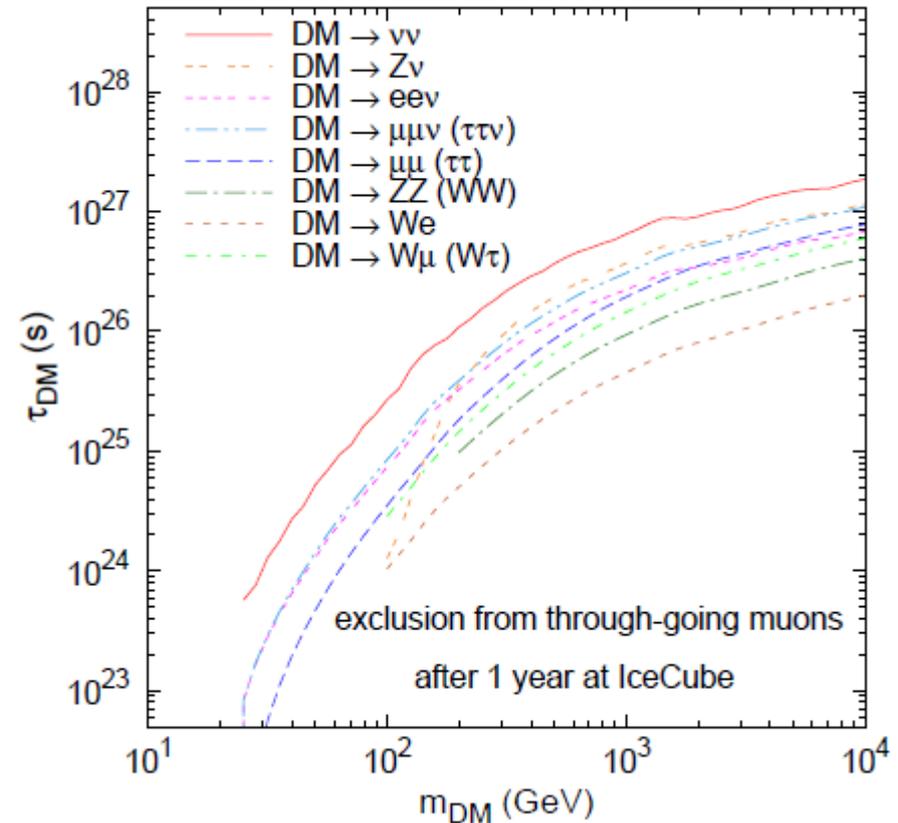
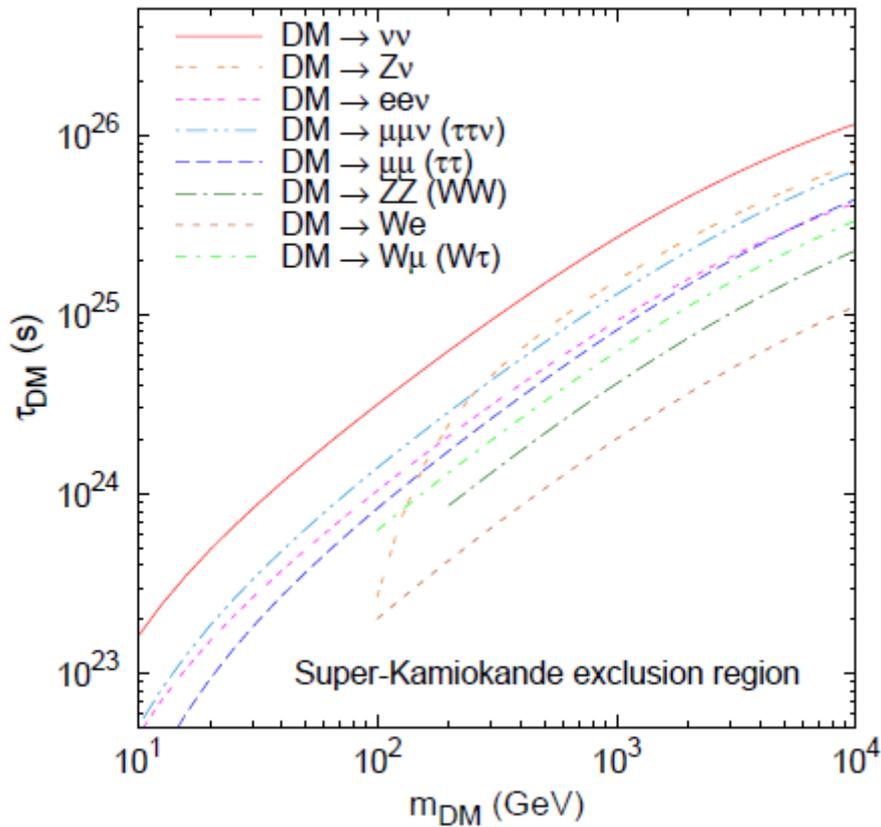
Covi, Grefe, AI, Tran  
arXiv:0912.3521

Typical lifetime of a fermionic dark matter particle,  
decaying via a dimension 6 operator suppressed by M

$$\tau_{\text{DM}} \sim 10^{26} \text{ s} \left( \frac{\text{TeV}}{m_{\text{DM}}} \right)^5 \left( \frac{M}{10^{15} \text{ GeV}} \right)^4$$

# Limits on the dark matter lifetime

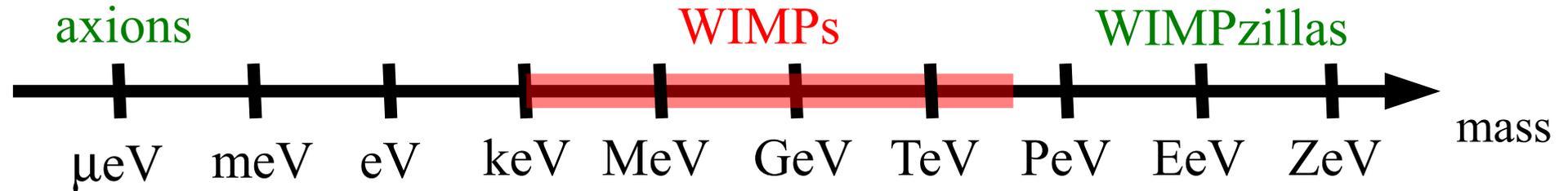
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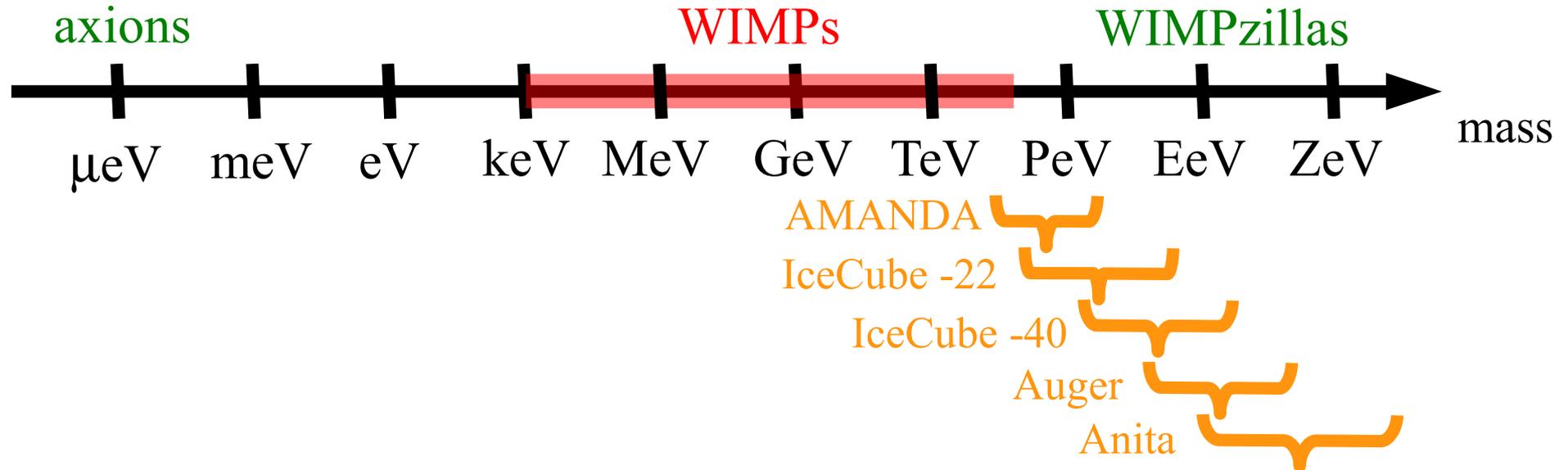
# Limits on the dark matter lifetime

Opening up the dark matter mass window...



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Opening up the dark matter mass window...



	$E_{\nu}^{\min} - E_{\nu}^{\max}$ (TeV)	$N_{\text{bg}}$	$N_{\text{sig}}$	$N_{\text{limit}}$
AMANDA	$16 - 2.5 \times 10^3$	6	7	5.4
IceCube-22	$340 - 2 \times 10^5$	0.6	3	6.1
IceCube-40	$2 \times 10^3 - 6.3 \times 10^6$	0.1	0	2.3
Auger	$10^5 - 10^8$	0	0	2.3
ANITA	$10^6 - 3.2 \times 10^{11}$	0.97	1	3.3

[arXiv:0705.1315](https://arxiv.org/abs/0705.1315)

[arXiv:1202.4564](https://arxiv.org/abs/1202.4564)

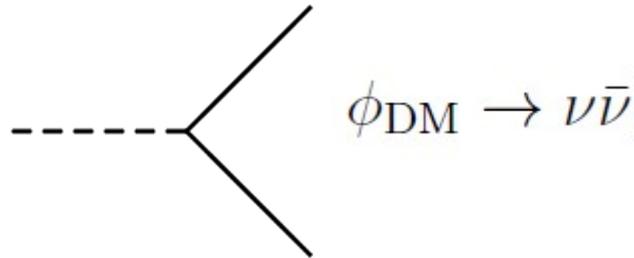
[arXiv:1103.4250](https://arxiv.org/abs/1103.4250)

[arXiv:1202.1493](https://arxiv.org/abs/1202.1493)

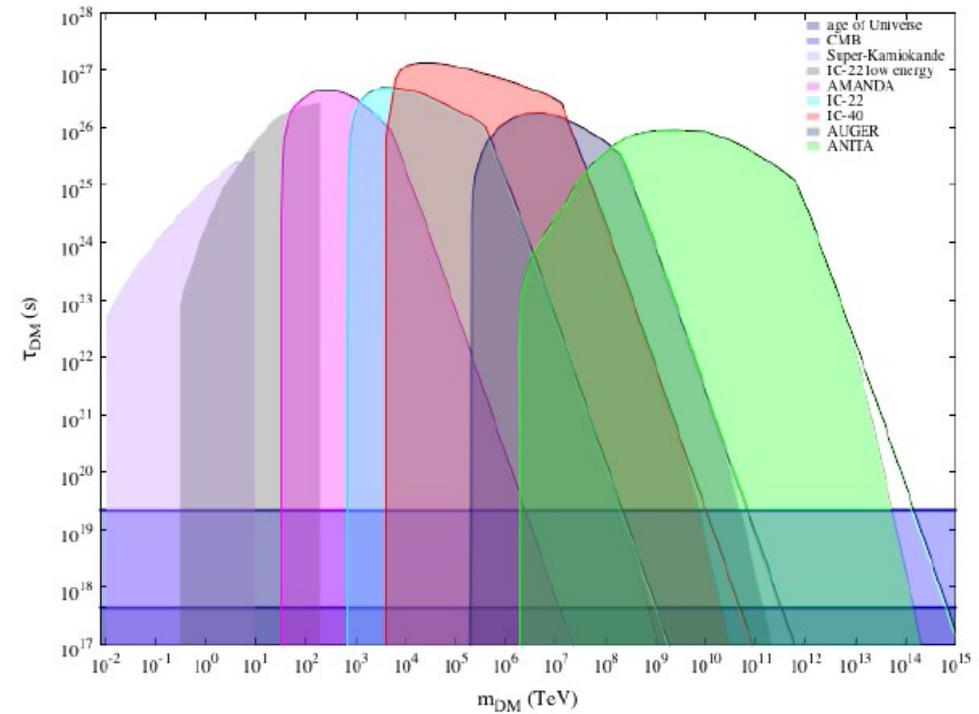
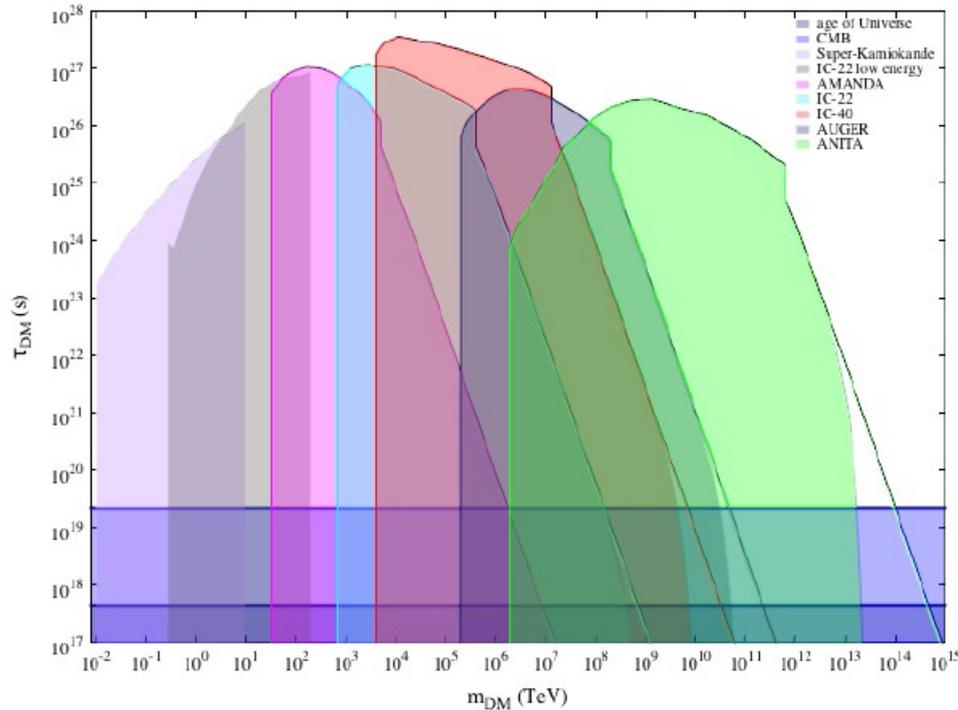
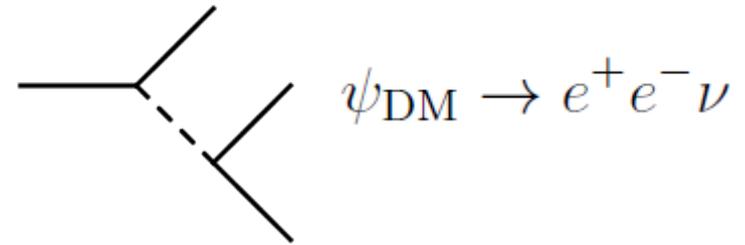
[arXiv:1011.5004](https://arxiv.org/abs/1011.5004)

# Limits on the dark matter lifetime

Monoenergetic neutrinos



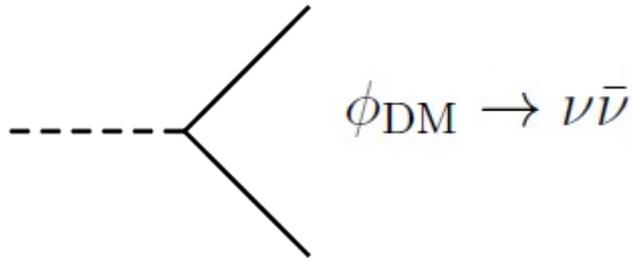
Continuous energy distribution



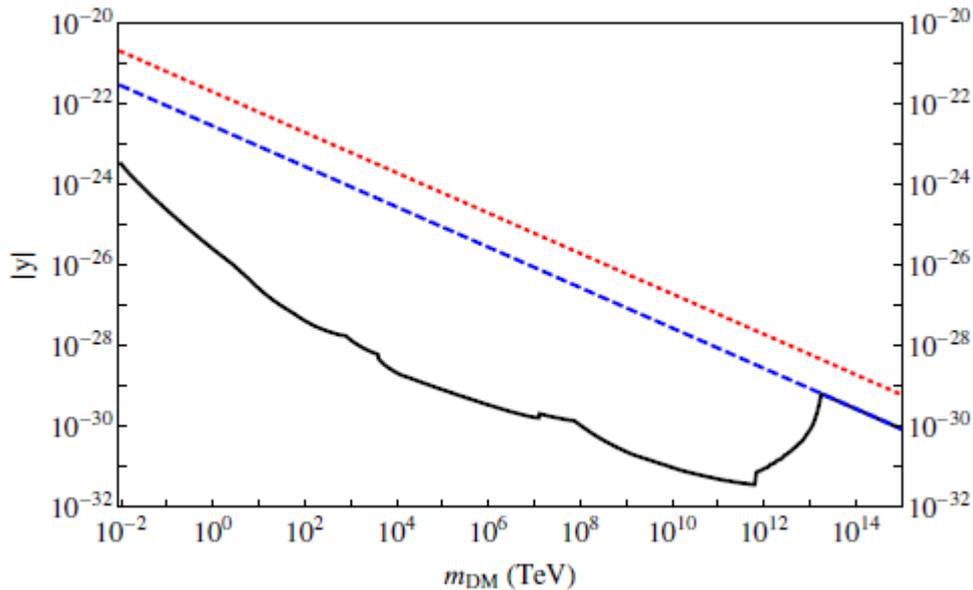
Esmaili, AI, Peres  
arXiv:1205.5281

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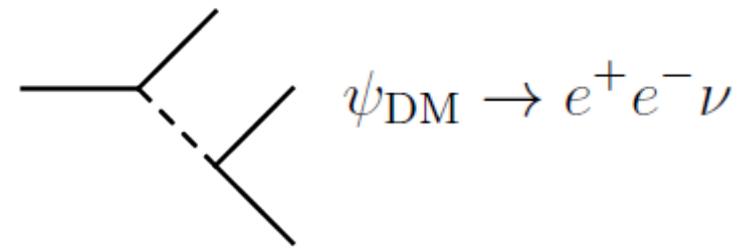
## Monoenergetic neutrinos



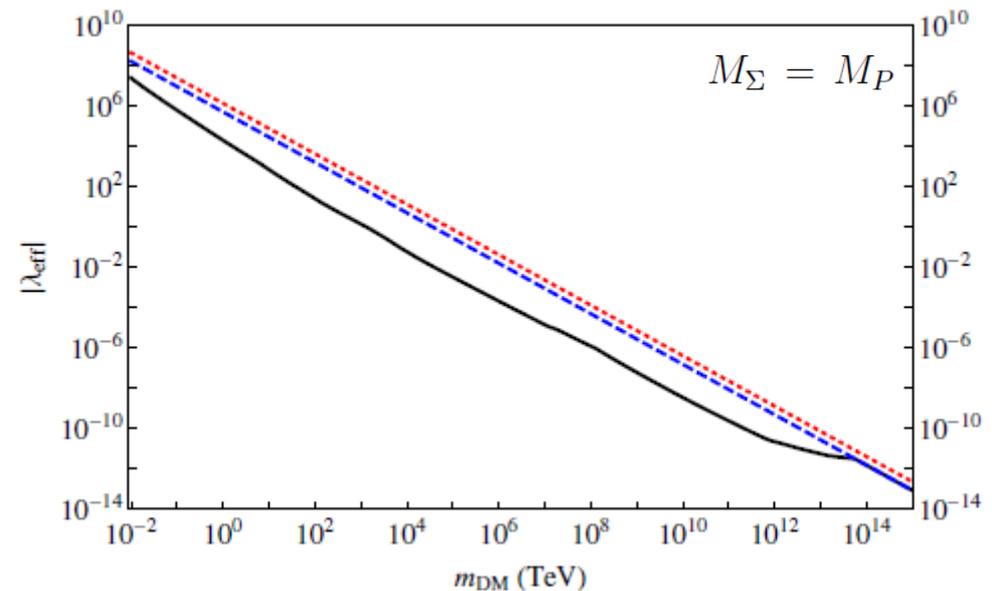
$$\Gamma(\phi_{\text{DM}} \rightarrow \nu\bar{\nu}) = \frac{|y|^2}{8\pi} m_{\text{DM}}$$



## Continuous energy distribution

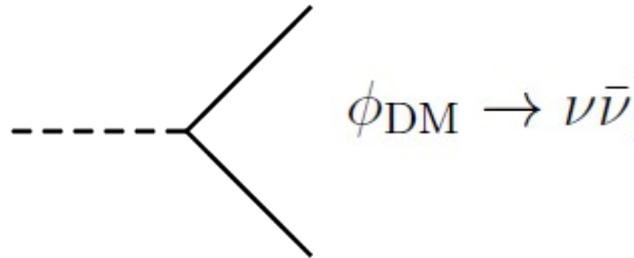


$$\Gamma(\psi_{\text{DM}} \rightarrow e^+e^-\nu) = \frac{|\lambda_{\text{eff}}|^4}{3072\pi^3} \frac{m_{\text{DM}}^5}{M_{\Sigma}^4}$$



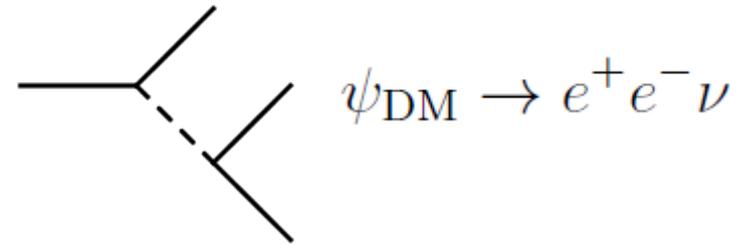
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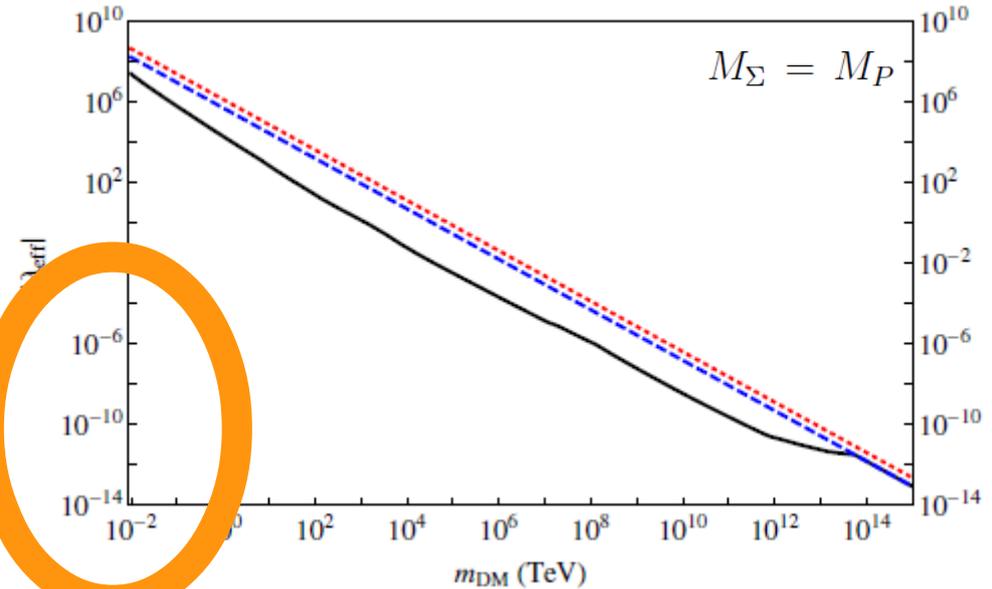
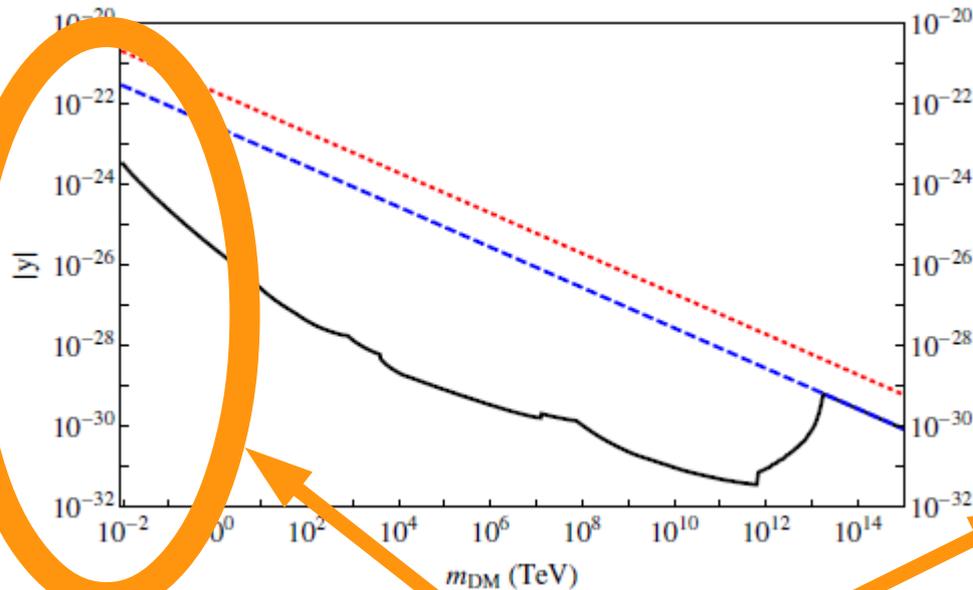


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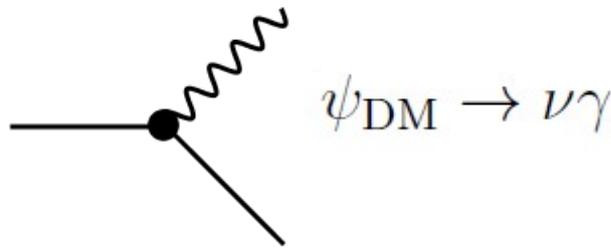
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*Symmetry?*

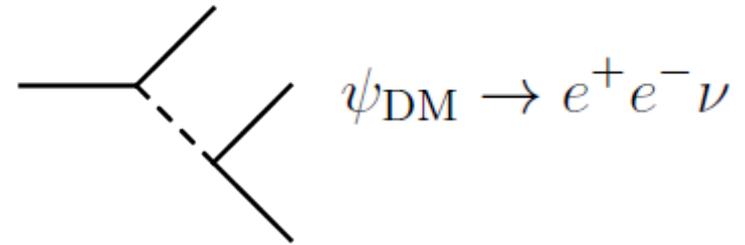
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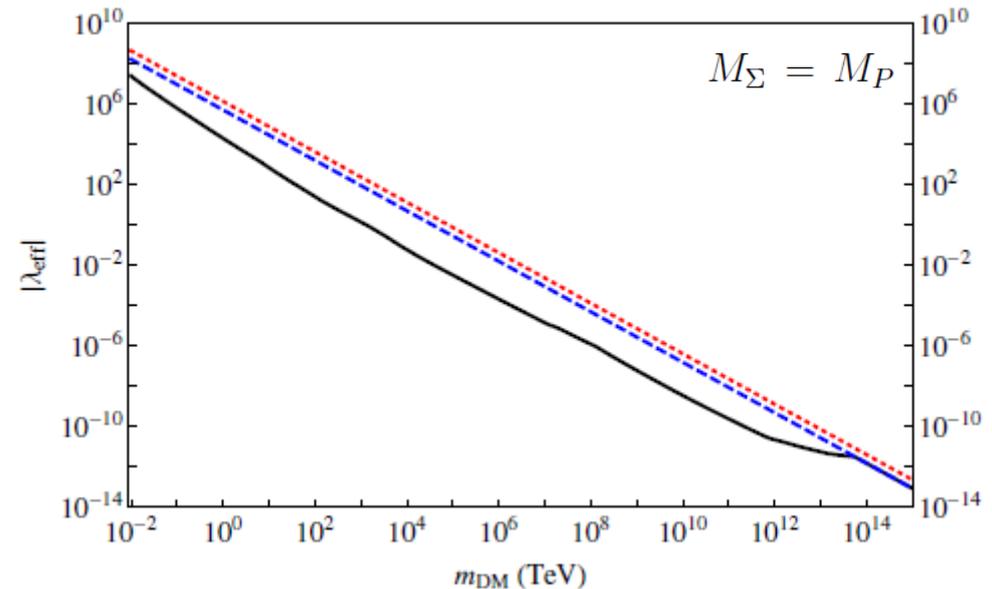
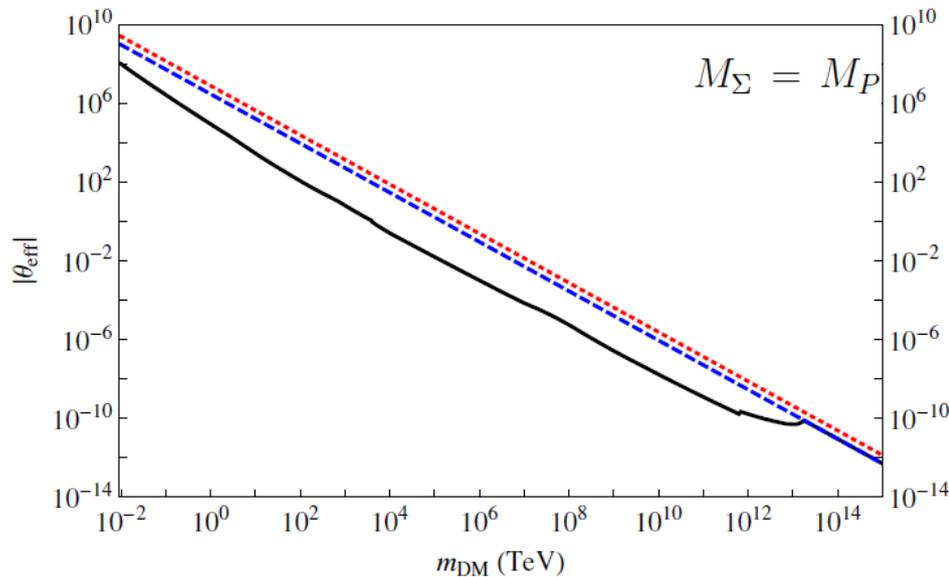


$$\Gamma(\psi_{\text{DM}} \rightarrow \nu \gamma) = \frac{|\mu_{\text{eff}}|^2}{8\pi} m_{\text{DM}}^3$$

## Continuous energy distribution

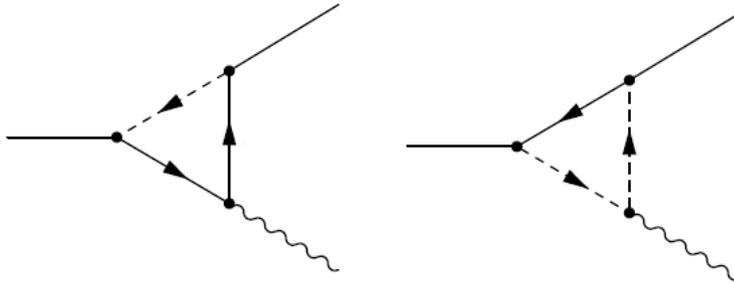


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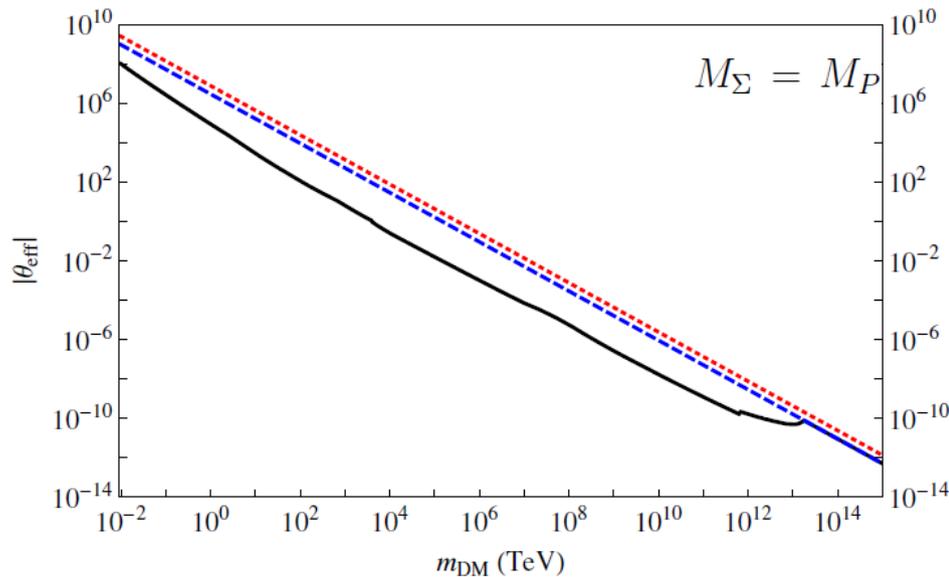
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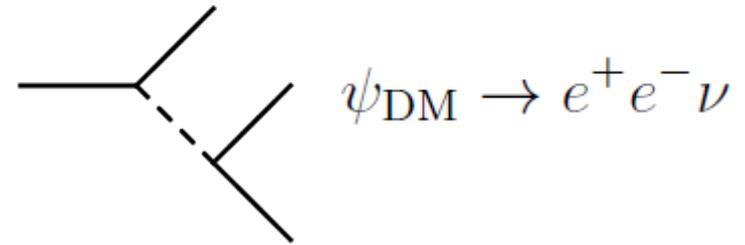


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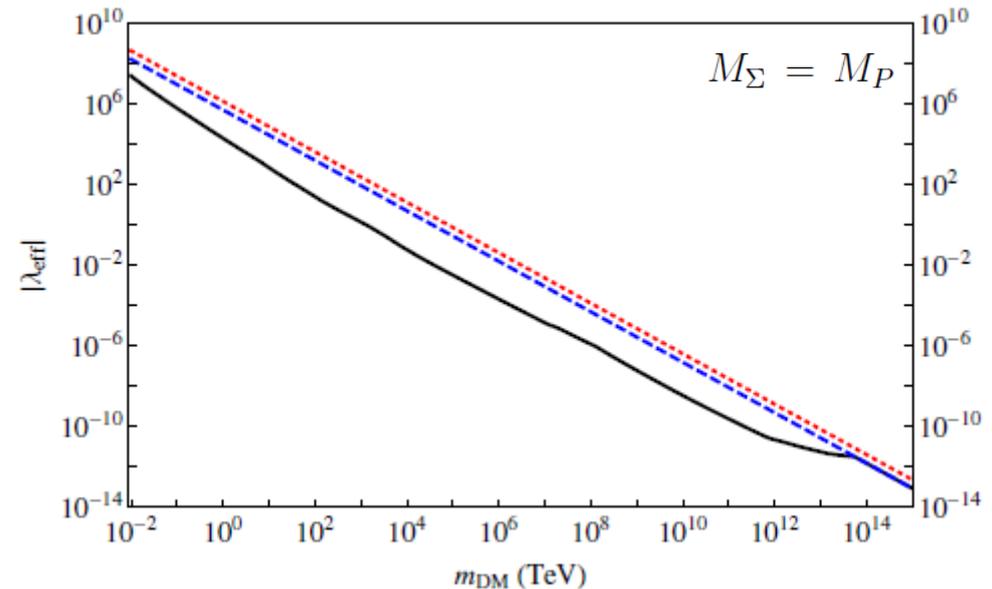
$$|\mu_{\text{eff}}| = \frac{e m_{\text{DM}} |\theta_{\text{eff}}|^2}{64\pi^2 M_{\Sigma}^2}$$



## Continuous energy distribution



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# Conclusions

Neutrinos are very elusive particles with tiny interaction cross sections with ordinary matter. Naively, a disadvantage when using neutrinos as messengers for indirect dark matter detection. In reality, an advantage (no deflection, no absorption), provided the small detection rates can be compensated by a large detector volume → SuperKamiokande, IceCube

- Limits on the self-annihilation cross section comparable, for large DM masses, to those from electrons/positrons.
- Best limits on the SD scattering cross section. Limits on the SI scattering cross section slightly worse than those from Xenon100, CDMS.
- Best limits on the lifetime of superheavy ( $m_{\text{DM}} > 10 \text{ TeV}$ ) DM particles.

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*Thank you for your attention!*